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# Unipolar nano-diode detector with improved performance using high-k material SiN<sub>x</sub>

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## Abstract

The performance of a solid-state planar nano-diode, namely self-switching diode (SSD), is improved by depositing a 100-nm-thick SiN<sub>x</sub> film with high dielectric constant (high-k) into the insulating nano-trenches. The SiN<sub>x</sub> film grown by using plasma-enhanced chemical vapour deposition (PECVD) can enhance the electric field coupling over the trenches and thus increase the accumulated charges for field effect. In this case, the current-voltage nonlinearity is improved significantly and the responsivity of high frequency rectification is also increased by a factor of almost one order up to 100 GHz. In addition, compared to the device without SiN<sub>x</sub> coating, the low frequency noise of the proposed diode is suppressed dramatically. The improved responsivity and noise-equivalent power (NEP) of 11 SSDs in parallel with SiN<sub>x</sub> coating are 110 V/W at 50 GHz and 180 pW/Hz<sup>1/2</sup>, which are comparable to the state-of-the-art data of reported SSD arrays.

Keywords: Self-switching diode, SiN<sub>x</sub>, High-k, Nano rectifier

## 1. Introduction

Terahertz technologies have been envisaged to have promising applications in medical and security imaging, pharmaceutical compound analysis, and radio astronomy etc. [1]. However, the frequency band is often called ‘terahertz gap’ due to a lack of high performance solid-state terahertz detectors and emitters at room temperature. Among the reported novel detectors [2, 3], a nano-scale self-switching diode (SSD) is a promising candidate for real applications. An SSD based on GaAs/AlGaAs heterostructure containing a two-dimensional electron gas (2DEG) has been shown to operate at frequencies up to 1.5 THz without any DC bias at room temperature [4]. Successful detection of free-space terahertz radiation between 1.3 and 2.5 THz at low temperatures from 10 to 150 K has also been demonstrated by SSDs fabricated using an InGaAs/InP 2DEG wafer [5]. In

addition, SSDs have been realized using other semiconductor materials, such as InAs/AlAs [6], GaN/AlGaN [7] and 2D materials [8]. For terahertz detectors, responsivity and noise performance are two most important factors. By connecting about 2000 SSDs in parallel, the thermal noise of the SSD array was significantly reduced and the noise-equivalent power (NEP) was improved to a value of 64 pW/Hz<sup>1/2</sup> [9], among the lowest ones of the reported terahertz detectors.

Additionally, modelling and simulation of SSD have been carried out intensively [10-13]. A number of simulations predicted that the SSD performance can be improved significantly if its nano-trenches are filled with high-k materials [14-17]. However, to the best of our knowledge, no fabricated device has been reported to demonstrate this yet. On the other hand, high-k materials have already been widely used as a substitute for SiO<sub>2</sub> in electronics to achieve high mobility, low threshold voltage and low leakage [18-21]. Examples are HfO<sub>2</sub> [22], Si<sub>3</sub>N<sub>4</sub> [23], SiO<sub>x</sub>N<sub>y</sub> [24] and

HfSiON [25]. In this work, SSD arrays are fabricated on 2DEG  $\text{In}_{53}\text{Ga}_{47}\text{As}$  substrate with high-k  $\text{SiN}_x$  filling for their nano-trenches, achieving high responsivity and low noise. Such improvement is highly desired for terahertz detection at room temperature.

## 2. Experiment

An  $\text{InGaAs/InAlAs}$  2DEG wafer from IQE Inc. was used to fabricate the SSDs. The 2DEG embedded in the  $\text{InGaAs/InAlAs}$  heterostructure is located 25 nm below the surface, as depicted in Figure 1(a). The carrier density and electron mobility of the 2DEG are  $1.3 \times 10^{12} \text{ cm}^{-2}$  and  $10,400 \text{ cm}^2/\text{V}\cdot\text{s}$  at 300 K, respectively. The relatively high electron mobility of the  $\text{InGaAs/InAlAs}$  based 2DEGs is beneficial for high frequency detection. The active area of the diode, consisting of asymmetric nanochannels defined with nano-trenches, was fabricated using electron-beam lithography, as shown in Figure 1(b). The SSD structure was then transferred into the substrate by bromine-based wet etching, which stops at the  $\text{InAlAs}$  buffer layer. The depth of the insulation trench is around 50 nm measured by atomic force microscopy (AFM), and the width is about 150 nm. The trenches can force the electrons to flow only through the nanochannels.

$\text{In}_{53}\text{Ga}_{47}\text{As}$	10 nm
$\text{In}_{52}\text{Al}_{48}\text{As}$	5 nm
N+ DD Si	$6.0 \times 10^{18}$
$\text{In}_{52}\text{Al}_{48}\text{As}$	10 nm
$\text{In}_{53}\text{Ga}_{47}\text{As}$	15 nm
$\text{In}_{52}\text{Al}_{48}\text{As}$	300 nm
Substrate	InP

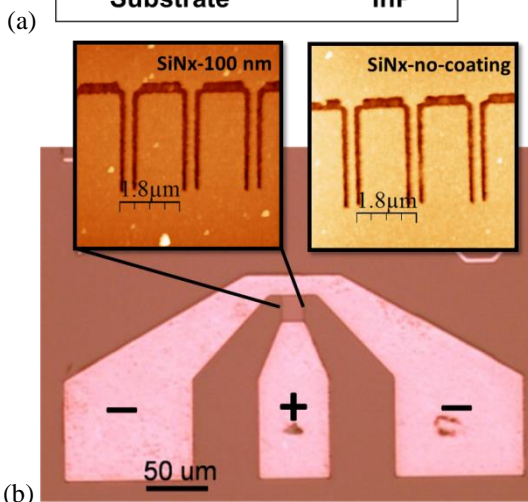


Figure 1. (a) Structure of  $\text{InGaAs/InAlAs}$  2DEG wafer. (b) AFM images of the etched SSD array with and without  $\text{SiN}_x$  coating with co-planar waveguide ohmic contacts. The SSD channel is 3  $\mu\text{m}$  long and 215 nm wide.

With the device unbiased, the effective channel is actually narrower than the geometric width due to the depletion region at the etched boundaries, which is introduced by the charges of surface states. The amount of surface charges could be affected by surface roughness, charge states, Fermi-level pinning and surface treatments [26]. The effective channel of the SSD can be widened or narrowed with forward or reverse bias respectively. This corresponds to the nonlinearity and thus a high-frequency rectification.

11 SSDs are accommodated in parallel on a mesa structure, as shown in Figure 1(b), which are fed by using a co-planar waveguide (CPW) with an impedance of 50  $\Omega$  for rectification characterisation. The signal and gap width of the CPW is 60  $\mu\text{m}$  and 55  $\mu\text{m}$ , respectively. To reduce the ohmic contact resistance, the electrodes are fabricated with  $\text{Au/Ge/Ni/Au}$ . The measured ohmic contact resistance is 0.86  $\Omega/\text{mm}$  and the specific resistance is  $4 \times 10^{-5} \Omega \cdot \text{cm}^2$ .

$\text{SiN}_x$  has been widely used as passivation for electronic devices, such as the field effect transistor (FET) [27] and high-electron-mobility transistor (HEMT) [28]. To improve the performance of the SSDs, the  $\text{SiN}_x$  coating was grown by plasma-enhanced chemical vapour deposition (PECVD), which offers excellent dielectric properties of the deposited and conformed step coverage [29]. The deposition was carried out with  $\text{SiH}_4$ ,  $\text{NH}_3$  and  $\text{N}_2$  at flow rates of 10, 16, 170 sccm, respectively, at 300  $^\circ\text{C}$ . The thickness of the  $\text{SiN}_x$  coating was 100 nm, namely twice that of the trench depth, so that the trenches were completely filled. The proposed device is nominated as ‘ $\text{SiN}_x$ -100 nm’ while another device without coating is nominated ‘ $\text{SiN}_x$ -no-coating’. These will be used for comparison in the following sections.

## 3. Results and discussion

### 3.1 I-V Characterisation

The devices were characterised at room temperature in the dark, by using Agilent E5270B semiconductor analyser. As the positive forward bias was increased, the current first increased rapidly and then quasi-linearly due to the series resistance, as shown in Figure 2(a). With a negative bias, the reverse current exists as the channel is not completely pinched off. Compared to that of the uncoated device, both the forward and reverse currents of the coated device are clearly increased, as shown in Figure 2(a). This is as expected from the enhanced field coupling over the trenches and the reduced electric field strength on the channel. However, note that the device nonlinearity, given by  $[I(V)+I(-V)]/2$ , is significantly enhanced particularly at low bias values, which are conditions desired for high-frequency rectification, as plotted in Figure 2(b). For instance, the nonlinearity of the current-voltage (I-V) curve is enhanced by a factor of about 18 at 0.1 V.

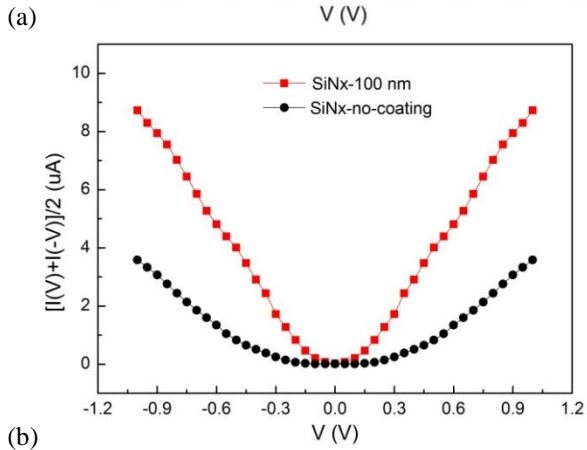
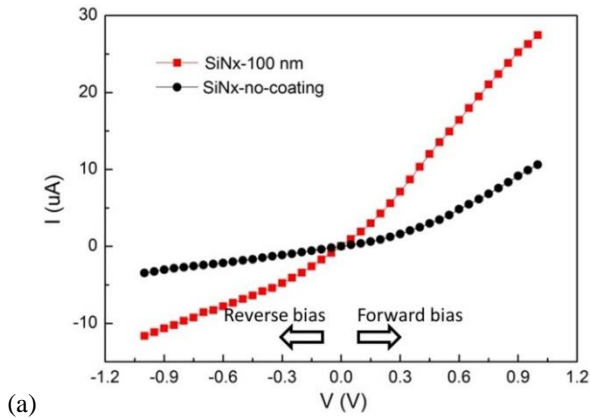


Figure 2. (a) The I-V characteristics of SiN<sub>x</sub>-100 nm and SiN<sub>x</sub>-no-coating. (b) The extracted nonlinear component of the I-V curves in (a).

### 3.2 RF Characterisation

An Agilent 8510XF vector network analyser (VNA) providing continuous frequency coverage from 45 MHz to 100 GHz was employed for high-frequency rectification measurement. The DC bias was provided with Keithley 2400 source measure unit through a bias-T, and the input power was calibrated before being fed into SSDs device. Figure 3(a) illustrates the rectified output voltages of the devices at zero bias. It can be seen that the output voltage of the device with SiN<sub>x</sub> coating is clearly higher than that of its uncoated counterpart, which reveals the strong enhancement of the responsivity. Figure 3(b) shows the output voltage as a function of input power measured at 50 GHz and zero bias. The output voltage is proportional to the input power, which corresponds to the square-law operation of the device. The extrinsic responsivity of the devices can be calculated by dividing the measured output voltage by the input power (-5.5 dBm). The extrinsic responsivity of the device SiN<sub>x</sub>-100 nm is 110 V/W at 50 GHz, which is much higher than that of SiN<sub>x</sub>-no-coating, 6.99 V/W. By integrating thousands of SSDs with high-k coating in parallel, the input impedance can further reduce from 33.3 kΩ (the zero-bias resistance) down to hundreds of ohms for better impedance matching and noise performance [9].

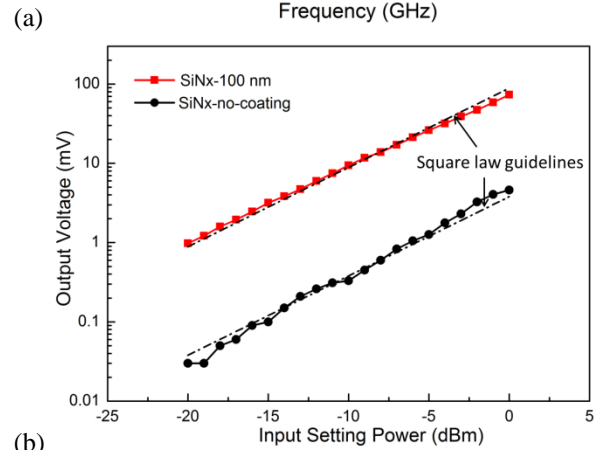
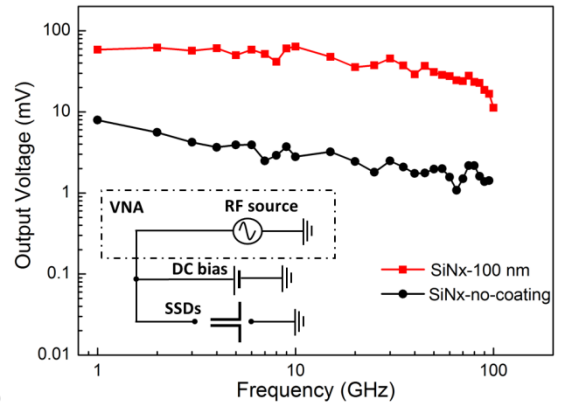


Figure 3. (a) Comparison of RF performance of SiN<sub>x</sub>-100 nm and SiN<sub>x</sub>-no-coating up to 100 GHz with an input power of -5.5 dBm at zero bias, and schematic representation of the RF measurement setup. (b) Output voltage as a function of input power for both devices measured at zero bias and 50 GHz.

### 3.3 Low-frequency Noise Measurement

The low-frequency noise performance of the devices was characterized by using a two-channel cross-correlation measurement [30, 31], as depicted in Figure 4(a). The two-channel cross-correlation method reduces the uncorrelated noise, but it cannot get rid of the input equivalent current noise of an operational amplifier (op-amp), which is mixed with the device noise in the noise measurement [30]. A low noise op-amp, TL1169, was then used to amplify the input noise signal due to its low input current noise of 1 fA/Hz<sup>1/2</sup> [32]. The highest frequency is approximately 10 kHz, which is limited by the amplifier bandwidth. The low-frequency noise is dominated by thermal and/or flicker noise [9, 33]. The measured noise spectral plot of the SiN<sub>x</sub>-100 nm and SiN<sub>x</sub>-no-coating at 32.5 nA bias and zero bias current is shown in Figure 4(b). It can be observed that the voltage noise significantly decreases by a factor of more than 10 times with the SiN<sub>x</sub> coating.

The voltage power spectral density is proportional to  $1/f^\beta$  at a certain current bias according to Equation 1 [9, 34].

$$S_V(f) = \frac{\alpha_H}{Nf\beta} I^2 R^2 \quad (1)$$

where  $I$  is the bias current,  $R$  is the resistance of the device under test,  $N$  is the total number of carriers in the SSD array,  $f$  is the frequency and  $\alpha_H$  is the Hooge's constant. The constant  $\beta$  is 0.96 by fitting the curve in Figure 4(b), which means the low frequency noise is dominated by  $1/f$  noise.

The dependence between the voltage noise,  $\sqrt{S_V(f)}$ , and the bias current at frequency of 100 Hz is shown in Figure 4(c). The voltage noise increases linearly as the bias current enlarges then the increase slows down. The linear dependency can be explained by Hooge's empirical equation, indicating that the low frequency noise is caused by the carrier mobility fluctuation [35]. At relatively high current bias, the effective channel width could increase slightly, leading to the slight decrease of the slope in Figure 4(c) according to Equation 2, which is expressed as

$$S_V(f) = \frac{1}{M^3} \frac{1}{n^3 q^2 \mu_n^2} \frac{L}{W_{eff} f^3} \frac{\alpha_H}{f\beta} I^2 \quad (2)$$

where  $n$  is the electron concentration,  $M$  is the number of SSDs in parallel,  $L$  is the channel length of the SSD,  $q$  is the electron charge, and  $\mu_n$  is the 2DEG electron mobility and  $W_{eff}$  is the effective channel width, which is 215 nm in this paper.

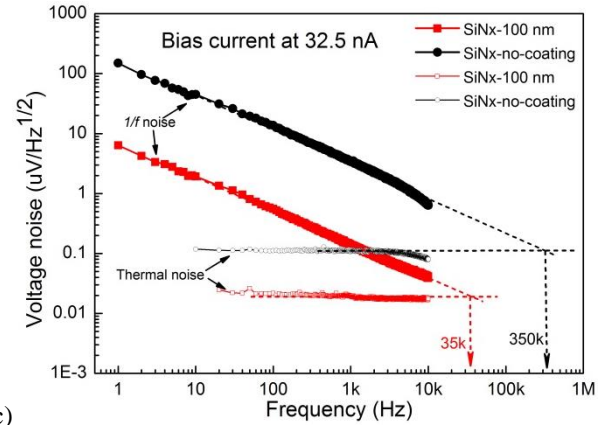
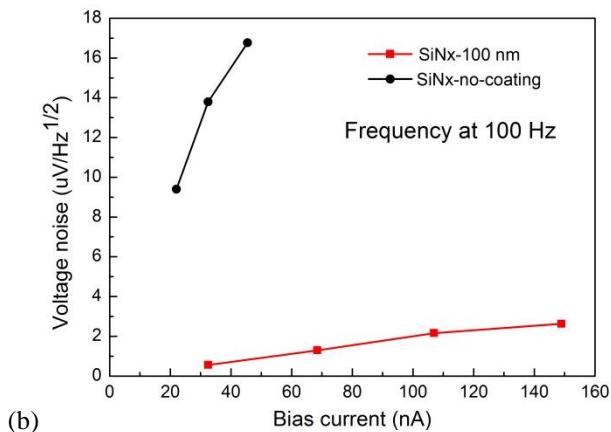
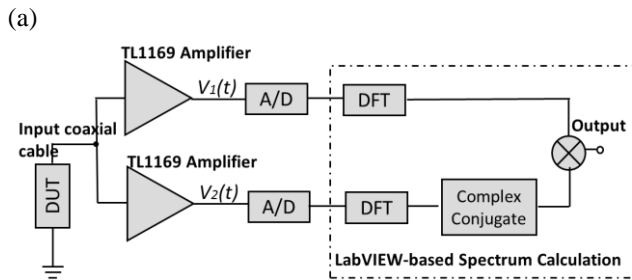


Figure 4. (a) Schematic system set-up of the two-channel cross-correlation measurement. (b) Voltage noise of device SiN<sub>x</sub>-100 nm and SiN<sub>x</sub>-no-coating at 32.5 nA and zero bias at room temperature with determination of the corner frequencies. (c) The linear dependence of voltage noise on the bias current at 100 Hz.

For the devices SiN<sub>x</sub>-100 nm and SiN<sub>x</sub>-no-coating, the channel length is 3 μm, and the zero-bias resistance is 33.3 and 616 kΩ, respectively. The device without coating has a Hooge's constant of  $3.7 \times 10^{-3}$ , which is lower than that of the device with coating ( $5.1 \times 10^{-2}$ ). Although the Hooge's constant of the device SiN<sub>x</sub>-100 nm is larger, the significant reduction of its resistance results in the reduced overall noise. In this case,  $1/f$  noise is caused by the scattering-induced carrier mobility fluctuations [36, 37]. The deposition of SiN<sub>x</sub> may introduce more traps at the interface in contact with the semiconductor. The increased interface states may cause more scattering, such as lattice scattering [38] and Coulomb scattering [39], resulting in a higher Hooge's constant. Using thermal evaporation may be able to decrease the traps and surface states caused by the PECVD coating process [40]. The corner frequency, shown as the cross point of the flick noise and thermal noise in Figure 4(a), which is 35 and 350 kHz for the device SiN<sub>x</sub>-100 nm and SiN<sub>x</sub>-no-coating at 32.5 nA bias current, respectively. This indicates that SiN<sub>x</sub> coating effectively suppresses the voltage noise and, thus, reduces the corner frequency. When the detected signal is smaller, the corner frequency could be further reduced for chopping [9]. NEP is another important metric of the sensitivity. When the frequency is above the corner frequency, thermal noise dominates the noise spectrum. Then the NEP is determined by thermal noise which is proportional to the resistance of SSDs. It is found that the extrinsic NEP significantly decreased by two orders of magnitude from  $1.6 \times 10^{-8}$  to  $1.8 \times 10^{-10}$  W/Hz<sup>1/2</sup> after SiN<sub>x</sub> coating. The NEP can be further reduced by integrating more SSDs due to the lower total resistance as well as the thermal noise. The responsivity and NEP of the 11 SSDs in parallel

are 110 V/W at 50 GHz and 180 pW/Hz<sup>1/2</sup> which are comparable to results achieved by integrating thousands of SSDs on an InGaAs/InAlAs 2DEG wafer [9].

#### 4. Conclusion

We proposed to use SiN<sub>x</sub> coating to improve the detection properties of the nano-rectifier, i.e. SSD. SSD arrays with nano-trenches coated with and without SiN<sub>x</sub> were fabricated, characterized, and compared. Those coated achieved high responsivity and low noise, agreeing with reported simulation results. They exhibit enhanced nonlinearity and output current, the parameters desired for room temperature THz detectors. This improvement is due to electric field coupling over the trenches causing an increase in the accumulated charges for the field effect. High-frequency measurement from 45 MHz to 100 GHz was carried out, which reveals a consistent improvement of the responsivity. The reduced low frequency noise, corner frequency and NEP demonstrate a good passivation of high-k SiN<sub>x</sub> for the SSDs.

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